# Dielectric Properties of Acceptor-Doped (Ba,Ca)(Ti,Zr)O<sub>3</sub> Ceramics

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Abstract. The dielectric properties of  $(Ba,Ca)(Ti,Zr)O_3$  ceramics containing various acceptor and donor dopants on the B sites have been studied. Formation of charge compensating "complexes" between acceptors, donors and oxygen vacancies has been observed. All acceptors cause a reduction of the Curie point. The maximum dielectric constant increases with the average grain size almost independent of the acceptor incorporated. Ceramics fired in reducive atmospheres reveal pronounced changes of the dielectric properties after annealing in oxidizing atmosphere. In particular the losses and the maximum dielectric constants are significantly affected by the annealing time, temperature and partial pressure of oxygen.

## 1. Introduction

Incorporation of heterovalent ions in the perovskite lattice of barium titanate leads to significant changes of the dielectric and ferroelectric properties. Generally, ions with large radius and low valency tend to enter the A (Ba) sites, while ions with small radius and higher valency favor the B (Ti) sites [1]. Ions bringing excess charges into their host lattice are named donor ions, while those bringing a charge deficit are named acceptor ions. Depending on their ionic radius, the trivalent ions of rare earths enter either A sites or B sites. A few of the smaller rare earth ions can enter both A and B sites, depending on the B/A atomic ratio of the perovskite [2]. These ions can therefore act as donors as well as acceptors in barium titanate.

Almost all heterovalent substituents tend to reduce the c/a ratio of the tetragonal unit cell and to decrease the ferroelectric Curie point. The presence of acceptors and donors strongly affects the stoichiometry and microstructure of ferroelectric perovskites [3]. The charge deficit of acceptors is usually compensated by doubly ionized oxygen vacancies,  $V_0^{\bullet\bullet}$  [4,5]. Formation of acceptor-oxygen vacancy complexes give rise to generation of elastic and electric dipols which have a strong influence on the dielectric and ferroelectric properties as well as the aging behavior [6].

This investigation deals with the influence of various acceptor-oxygen vacancy and acceptordonor complexes on the ferroelectric and dielectric properties of reductively fired and oxidizingly aftertreated  $(Ba,Ca)(Ti,Zr)O_3$  (BCTZ). This basic composition has been selected because of its importance for the manufacture of ceramic capacitors and the strong effect of re-oxidation treatment on the dielectric properties.

### 2. Sample Preparation

Ceramic materials of composition  $(Ba_{0.96}Ca_{0.04})$  $[Ti_{1-x-y-z}A_xD_yZr_z]_aO_{3-\delta}$  were prepared by standard solid state techniques, using reagent grade raw materials. As acceptors the ions of A = Mg, Al, Ga, Mn, Cr, Sc, Y, Yb have been studied and as donors D = Nb and W. The acceptor and donor concentrations were x, y  $\leq 0.01$ , while that of Zr was z = 0.18and a = 0.995.

Before and after calcination the powders were brought into a fine grained state with particle sizes between 0.5 and 1  $\mu$ m, using a ball mill with 2 mm Ø Y-stabilized zirconia balls. The major impurities Sr in  $BaCO_3$  and Hf in  $ZrO_2$  were taken into account at weighing. After 6 h calcination at  $1250^{\circ}C$  in air single- phase perovskites have been observed in X-ray diffraction analyses.

The fine milled calcined powders were granulated and uniaxially pressed at 4 kbar to green discs of about 0.6 mm thickness and 6 mm diameter. The excess of A-site ions, B/A < 1, required small SiO<sub>2</sub> additives as sintering aid. The predensified green discs were fired for 1 h at 1320°C in a reductive atmosphere of moist (saturated with H<sub>2</sub>O at 20°C) N<sub>2</sub>, containing 1 vol.% H<sub>2</sub>. The partial pressure of oxygen was continuously monitored, using a zirconia cell at 735°C (AMETEK, Pittsburgh). The measured values of Po<sub>2</sub> agreed well with the calculated  $Po_2 = 2.1 \cdot 10^{-5}$  Pa. Sintered ceramic samples exhibited densities of 95% to 99% of the theoretical density. Part of the ceramics were aftertreated for 2 h at 1000°C in air  $(Po_2 = 2 \cdot 10^4 Pa)$  or in an atmosphere of N<sub>2</sub>/50 ppm  $O_2$  (Po<sub>2</sub> = 5 Pa). Heating and cooling rates were 300°C/h or 600°C/h.

The average grain sizes  $d_{av}$  have been determined from etched ceramics, using the line intercept method and light microscopy. The grain size decreased with increasing acceptor dope concentration. A detailed list of the measured grain sizes is shown in Table 1. Cr/Ni-Au electrodes were applied by evaporation on the sintered ceramic discs. Curie point and dielectric properties were determined at 1 kHz and 1 Vrms in the temperature range  $-50^{\circ}$ C to  $150^{\circ}$ C. The Curie point,  $T_{\rm C}$ , was determined from the maximum of the dielectric constant  $\varepsilon$  in the dielectric temperature characteristic  $\varepsilon(T)$ . All dielectric samples prepared insulation showed very high resistances  $R_{\rm I} > 10^{12} \Omega {\rm cm}$  at 20°C.

## 3. Experimental Results

#### 3.1. Curie Point Shifts

The Curie point  $T_{\rm C}$  sensitively depends on the kind of acceptor ion on B sites. The  $T_{\rm C}$  shift is determined by the ionic radius as well as the valence state and the ground state of the acceptor ion [7]. Both the acceptor ion and the presence of oxygen vacancies tend to remove locally the tetragonal symmetry and the ferroelectric state, thus leading to a reduction of  $T_{\rm C}$ . The presence of the small Ca content seems not to affect the acceptor concentration for B/A ratios

*Table 1.* Average grain size and maximum permittivity of reducively fired (1320°C,  $Po_2 = 2.1 \cdot 10^{-5} Pa$ ) and re-oxidized (1000°C,  $Po_2 = 5 Pa$ ) ( $Ba_{0.96}Ca_{0.04}$ )[ $Ti_{0.82-x}A_xZr_{0.18}$ ]<sub>0.995</sub> $O_{3-\delta}$ .

Acceptor (A)	x	$d_{av}$ ( $\mu m$ )	$\varepsilon_{\rm max}$ reduced	ɛ <sub>max</sub> re−oxid
Mg	0.0025	5.7	15,900	26,250
	0.0050	4.7	11,200	16,700
	0.0075	3.6	9,160	13,610
	0.0100	2.4	6,750	8,800
Al	0.0025	4.4	11,600	16,080
	0.0050	3.3	10,350	15,060
	0.0075	2.8	7,590	10,310
	0.0100	2.6	8,150	10,430
Cr	0.0025	6.1	12,790	26,060
	0.0050	5.4	13,050	15,990
	0.0075	5.0	11,350	12,540
	0.0100	4.1	9,320	9,610
Mn	0.0025	2.9	11,860	13,250
	0.0050	2.6	8,940	10,360
	0.0075	2.5	8,660	7,231
	0.0100	2.4	7,230	7,120
Sc	0.0025	6.7	15,610	28,090
	0.0050	5.9	13,690	24,050
	0.0075	4.6	10,550	18,670
	0.0100	3.9	9,800	15,800
Y	0.0025	7.0	17,390	28,650
	0.0050	5.0	13,180	20,610
	0.0075	3.9	11,120	15,220
	0.0100	3.4	9,970	13,120
Yb	0.0025	7.2	16,950	29,520
	0.0050	5.6	13,140	22,020
	0.0075	4.1	10,140	15,980
	0.0100	3.3	8,910	14,370

 $\simeq 0.995$ . Also no Curie point change was obtained for Ca contents < 8 at.%, while  $(Ca_{Ti})''$  acceptors strongly reduce  $T_C$  [8]. Co-doping with even small amounts of  $(Nb_{Ti})$  donors leads to semiconductive ceramics after firing in reducing atmosphere. Since the Nb donors have not been compensated by Ca acceptors, we can assume that  $(Ca_{Ti})''$  are not present.

Figures 1(a) and 1(b) show the variation of  $T_{\rm C}$  as function of the dopant concentration x for various acceptor ions A = Mg<sup>2+</sup>, Al<sup>3+</sup>, Ga<sup>3+</sup>, Sc<sup>3+</sup>, Y<sup>3+</sup>, Yb<sup>3+</sup>, Mn<sup>n+</sup> and Cr<sup>n+</sup> in the perovskite phase (Ba<sub>0.96</sub>Ca<sub>0.04</sub>)[Ti<sub>0.82-x</sub>Zr<sub>0.18</sub>A<sub>x</sub>]<sub>0.995</sub>O<sub>3- $\delta$ </sub>. The solid lines in Fig. 1 represent the materials sintered in reductive atmosphere (Po<sub>2</sub> = 2·10<sup>-5</sup> Pa), while the dashed lines represent the same materials aftertreated at Po<sub>2</sub> = 5 Pa. The atomic ratio of A and B site ions was adjusted to B/A = 0.995. Corresponding to B/A < 1 one can be sure that all dopant ions act as acceptors on B sites, i.e., (Mg<sub>Ti</sub>)", (Ga<sub>Ti</sub>)', (Sc<sub>Ti</sub>)', (Y<sub>Ti</sub>)', (Yb<sub>Ti</sub>)'.



*Fig.* l(a), l(b). Concentration dependence of the Curie point for various acceptor-doped materials of composition  $(Ba_{0.96}Ca_{0.04})[Ti_{0.82-x}Zr_{0.18}A_x]_{0.995}O_{3-\delta}$  in the reduced  $(Po_2 = 2 \cdot 10^{-5} Pa)$  and oxidized  $(Po_2 = 5 Pa)$  state, represented by the solid and dashed lines, respectively. The dashed-dotted lines for  $Mn^{3+}-Mn^{4+}$  and  $Cr^{4+}$  represent ceramics oxidized in air  $(Po_2 = 2 \cdot 10^4 Pa)$ .

Former investigations revealed that  $Mn^{2+}$ , resp.  $(Mn_{Ti})''$  is formed in the given reductive atmosphere [4].  $Cr^{3+}$  ions form  $(Cr_{Ti})'$ . The rare earth doped materials exhibit a high insulation resistance and a large  $T_C$  shift after reductive firing, thus indicating that Y and Yb have mainly entered the B sites. However, small amounts of rare earth ions on A sites cannot be excluded.

As already expected, all acceptors caused a negative Curie point shift,  $\Delta T_{\rm C} = [T_{\rm C}(x) - T_{\rm C}(0)] < 0$ . The largest shifts occurred for the divalent ions, yielding for

$$Mn^{2+}, Mg^{2+}: \Delta T_{\rm C}(A^{2+}_{\rm (Ti)})$$
  
\$\approx -80°C/at.\% in BCTZ-18\%Zr

This shift compares well with that observed for  $Ca^{2+}$  on Ti sites [8]. The shift of the trivalent ions is much lower.

$$Cr^{3+}, Sc^{3+}, Y^{3+}, Yb^{3+}: \Delta T_{\rm C}(A^{3+}_{({\rm Ti})})$$
  
 $\simeq -50^{\circ}{\rm C/at.\%}$  in BCTZ-18%Zr

For comparison, the Curie point shift for isovalent  $Zr^{4+}$  ions is (dotted line in Fig. 1(a)):

$$\Delta T_C(Zr_{(Ti)}^{4+}) = -7^{\circ}C/at.\%$$

The large difference in the  $T_{\rm C}$ -shift between ions of different valencies already reflects the important role of the oxygen vacancies on the ferroelectric state. In the case of isovalent Zr<sup>4+</sup> no oxygen vacancies are formed, so that the ferroelectric state is weakly disturbed and the resulting reduction of  $T_{\rm C}$  is small.

In the case of trivalent substituents, one  $V_0^{\bullet}$  is formed for charge compensation of 2 acceptors. Divalent substituents acting as twofold charged acceptors on B sites need one  $V_0^{\bullet}$  for charge compensation. Divalent acceptor ions thus give rise to the largest  $\Delta T_C$ . The presence of one oxygen vacancy together with a lower valent ion on B sites obviously causes a local reduction of c/a in a much larger volume as compared to  $Zr^{4+}$ . Similar results were obtained for compositions with a lower Zr content (y = 0.10). However, in this case the  $T_C$ -shifts were slightly lower.

3.1.1. Re-Oxidation Effects on  $T_{\rm C}$ . The annealing procedure leads to a shift  $\delta T_{\rm C} = (T_{\rm C(ox)} - T_{\rm C(red)})$  of the Curie point. This  $T_{\rm C}$ -shift between the reduced and oxidized state is usually smaller than the compositional shift  $\Delta T_{\rm C}$ . The Curie point of the oxidized state  $T_{\rm C}({\rm ox})$  was determined in samples annealed for 2 h at  $1000^{\circ}$ C in N<sub>2</sub>/50 ppm O<sub>2</sub> (Po<sub>2</sub> = 5 Pa) (heating and cooling rate: 300°C/h). In the case of valence stable acceptors, e.g., Mg, Y, Sc, Yb the effect of reoxidation on  $T_{\rm C}$  is expected to be rather small, because the acceptor  $V_0^{\bullet\bullet}$  complex should be not affected by the oxidizing heat treatment, except for local rearrangements of oxygen vacancies. A very small  $\delta T_{\rm C}$  was observed for Mg and Y doped material, as can be seen in Fig. 1. In the case of Sc and Yb doped BCTZ a slightly negative Curie point shift of  $\delta T_{\rm C} = -8^{\circ}$ C/at.% Yb and  $-13^{\circ}$ C/at.% Sc has been found.

In contrast to the negative  $\delta T_{\rm C}$  of ions with fixed valency the Curie point shift turns out to be positive

for Mn and Cr acceptors. This effect is expected because the number of oxygen vacancies is reduced, due to the oxidation of these valence instable ions. TGA analyses have confirmed that under the conditions of reducive firing (Po<sub>2</sub> =  $2.1 \cdot 10^{-5}$  Pa, 1320°C) Mn<sup>2+</sup> is formed, while at 1000°C in an atmosphere of N<sub>2</sub>/50 ppm O<sub>2</sub> (Po<sub>2</sub> = 5 Pa)  $Mn^{3+}$  is the prevailing valence state, as can be seen in Fig. 2. At even higher partial pressure of oxygen and lower temperature, e.g.,  $800^{\circ}$ C—air (2·10<sup>4</sup> Pa), about 50% of the  $Mn^{3+}$  are transformed to  $Mn^{4+}$  [4]. This leads to a lower  $\Delta T_{\rm C}$  as compared to the Zr-doped material (line marked with  $Mn^{3+}-Mn^{4+}$  in Fig. 1(a)). In contrast to Mn<sup>2+</sup> the compositional Curie point shift of Mn<sup>3+</sup> is rather small,  $\Delta T_{\rm C}({\rm Mn}^{3+}) \simeq -15^{\circ}{\rm C/at.\%}$ . This effect is expected, because part of the strained volume is restored with respect to higher c/a ratio, giving rise to a higher  $T_{\rm C}$ . The oxidation of  ${\rm Cr}^{3+}$  leads to comparable results. The dashed lines in Fig. 1(b) marked with  $Cr^{3+}-Cr^{4+}$  and  $Cr^{4+}$  represent the ceramics annealed in 50 ppm  $O_2$  (Po<sub>2</sub> = 5 Pa) and in air ( $Po_2 = 2 \cdot 10^4 Pa$ ), respectively.

Generally, the level of  $T_{\rm C}$  depends on the annealing temperature  $T_{\rm a}$ , the annealing time  $t_{\rm a}$  and the partial pressure of oxygen. The influence of Po<sub>2</sub> and  $T_{\rm a}$  was studied for BCTZ containing Mg, Sc, Cr, Mn and one additional sample containing Mn/W. These materials were annealed for  $t_{\rm a} = 2$  h at various top temperatures  $T_{\rm a}$  in N<sub>2</sub>/50 ppm O<sub>2</sub> or air. Heating and cooling rates were 600°C/h. The results are shown for N<sub>2</sub>/50 ppm O<sub>2</sub> (dashed lines) and air (solid lines) in Fig. 3. Ions with fixed valency (Mg<sup>2+</sup>, Sc<sup>3+</sup>) exhibit the same



*Fig.* 2. Oxygen consumption of  $(Ba_{0.96}Ca_{0.04})$ -[Ti<sub>0.81</sub>Zr<sub>0.18</sub>Mn<sub>0.01</sub>]<sub>0.995</sub>O<sub>3- $\delta$ </sub>, due to valence change of Mn. TG equilibrium measurement as function of the oxygen pressure p<sub>O2</sub> various temperatures.



*Fig. 3.* Variation of the Curie point with the annealing temperature for different acceptor doped materials of composition  $(Ba_{0.96}Ca_{0.04})[Ti_{0.815}Zr_{0.18}A_{0.005}]_{0.995}O_{3-\delta}$  with A = Cr, Mn, Mg, Sc. The lines marked with Mn-W represent a material of composition  $(Ba_{0.96}Ca_{0.04})[Ti_{0.802}Zr_{0.178}Mn_{0.01}W_{0.0075}]O_{3-\delta}$ . The materials were annealed for 2 h in an oxidizing atmosphere with Po<sub>2</sub> = 5 Pa (dashed lines) and Po<sub>2</sub> = 2·10<sup>4</sup> Pa (solid lines).

behavior, a slight decrease of  $T_{\rm C}$  at annealing temperatures  $T_{\rm a}$  ranging between  $\simeq 500^{\circ}{\rm C}$  and 700°C. The influence of Po<sub>2</sub> on  $T_C$  (N<sub>2</sub>/50 ppm O<sub>2</sub> or air) is rather low. The valency changes of Mn<sup>2+</sup> to  $Mn^{3+}/Mn^{4+}$  and  $Cr^{3+}$  to  $Cr^{4+}$  both give rise to an increase of  $T_{\rm C}$  at higher  $T_{\rm a}$  and Po<sub>2</sub>. For  $T_{\rm a} \leq 700^{\circ} {\rm C}$ and N<sub>2</sub>/50 ppm O<sub>2</sub> the  $T_{\rm C}$  shift of Cr<sup>3+</sup> behaves similar to the ions with a fixed valency, while at higher  $T_{\rm a}$  the increase of  $T_{\rm C}$  indicates the onset of the oxidation to  $Cr^{4+}$ . The valence change of  $Mn^{2+}$ gradually takes place already at lower annealing temperatures. For annealing in air the data of Cr- and Mn-doped BCTZ reveal a saturation of  $T_{\rm C}$  at high  $T_{\rm a}$ . From these results can be derived that the oxidation of the majority of the incorporated Cr and Mn ions only can be reached at  $T_{\rm a} > 800^{\circ}$ C and high oxygen pressure. The material doped with 1 at.% Mn and 0.75 at.% W (Mn-W) exhibits in contrast to single Mn-doped BCTZ a very small  $T_{\rm C}$ -shift vs  $T_{\rm a}$  for both atmospheres. The small Curie point shift observed confirms the formation of stable Mn-W complexes which cannot be oxidized at all.

Due to the high Zr content all materials investigated exhibit a diffuse phase transition [9,10]. Such materials are characterized by a broad dielectric maximum as a function of temperature and a negative value  $(T_{\rm C} - T_0)$ , where  $T_0$  denotes the Curie temperature. With increasing acceptor concentration the diffuse character of the phase transition is enhanced. Also, increasing ionic radius of the acceptors enhances the negative value of  $(T_{\rm C} - T_0)$ , due to an increased distortion of the lattice. Figure 4 represents  $(T_{\rm C} - T_0)$  of  $(\text{Ba}_{0.96}\text{Ca}_{0.04})$ - $[\text{Ti}_{0.82-x}\text{Zr}_{0.18}\text{A}_x]_z\text{O}_{3-\delta}$  for a constant concentration x = 0.0075 of various acceptors A in the reduced state (solid line, solid symbols) and the oxidized state (dashed line, open symbols). The diffuse character is even stronger for divalent acceptors (Mg<sup>2+</sup>) as compared to trivalent acceptors with comparable ionic radius. It should be noted that the diffuse character of the phase transition is intensified by oxidation.

3.1.2. Donor/Acceptor Complexes. In the case of mixed donor/acceptor doped BCTZ parts of the acceptors are compensated by donors, due to formation of donor-acceptor "complexes". The term complex implies the presence of an average electronically compensated donor-acceptor unit with varying donor-acceptor distance, corresponding to the statistical distribution of the ions at high temperatures. Such materials require less amounts of  $V_0^{\bullet\bullet}$  for charge compensation of the acceptors. Although Nb donors usually give rise to a strong depression of  $T_C$  in BaTiO<sub>3</sub> [11], the increase of the Nb concentration leads to an increase of  $T_C$  in reducively fired Nb-Mn doped BCTZ. In complex Mn-Nb and Mn-W doped



*Fig.* 4. Difference of the Curie point and the Curie temperature as a function of the ionic radius for different acceptor doped materials of composition  $(Ba_{0.96}Ca_{0.04})$ - $[Ti_{0.8125}Zr_{0.18}A_{0.0075}]_{0.995}O_{3-\delta}$  in the reduced  $(Po_2 = 2 \cdot 10^{-5} Pa)$  and oxidized  $(Po_2 = 5 Pa)$  state, represented by the solid symbols (solid line) and open symbols (dashed line), respectively.

materials elimination of  $V_0^{\bullet\bullet}$  obviously gives rise to the observed increase of  $T_C$ . The positive effect of  $V_0^{\bullet\bullet}$ elimination on  $T_C$  thus over-compensates the depressive effect of Nb and W donors on B sites.

Figure 5 demonstrates increasing Curie points for increasing donor concentration x in reduced BCTZ and decreasing Curie points in oxidized (Po<sub>2</sub> =  $2 \cdot 10^4$  Pa) BCTZ of composition

$$(Ba_{0.96}Ca_{0.04})[Ti_{0.81-x}Zr_{18}Mn_{0.01}D_x]_aO_{3-\delta}$$
  
$$D = (Nb_{Ti}^{5+}), (W_{Ti}^{6+})''$$

The difference in Curie point between reduced and oxidized BCTZ becomes smaller and smaller with increasing donor concentration x. At a Nb concentration of  $x \simeq 0.01$  and a W<sup>6+</sup> concentration of  $x \simeq 0.05$ the Curie points of reduced and reoxidized material have widely approached each other, i.e.,  $\delta T_{\rm C} \simeq 0$ . From this can be deduced that for given atomic ratios Mn/Nb = 1:1 and Mn/W = 2:1 the charges of donors and acceptors are compensating each other. Thus, in the case of Mn/Nb doped material the charge complex  $[Mn'-Nb^{\bullet}]^{0}$  and for Mn/W  $[2Mn'-W^{\bullet}]^{0}$  are formed. These complexes are stable in reducive ( $\simeq 10^{-6}$  Pa) as well as in oxidizing atmosphere  $(2 \cdot 0^4 \text{ Pa})$ . From Fig. 5 it is apparent that the oxidation of Mn applies primarily to the Mn which is not involved in the formation of charge complexes. The crossing points



*Fig.* 5. Variation of the Curie point with increasing donor content for materials of composition  $(Ba_{0.96}Ca_{0.04})$ - $[Ti_{0.81-x}Zr_{0.18}Mn_{0.01}D_x]_{3-\delta}$  with D = Nb, W. The results in the reduced  $(Po_2 = 2 \cdot 10^{-5} Pa)$  and oxidized  $(Po_2 = 5 Pa)$  state are represented by solid symbols (solid lines) and open symbols (dashed lines), respectively. The lines marked with  $Mn^{2+}$  (x = 0) and  $Mn^{3+}$  (x = 0) represent  $T_C$  for the corresponding material without donors.

(arrows in Fig. 5) of the  $T_{\rm C}$  curves for the reduced and oxidized state slightly deviate from the expected values, e.g., Mn : Nb = 1 : 1, Mn : W = 2 : 1. This may be attributed to the disordered distribution of donors and acceptors in the material. In the case of Mn/W-doped materials also charge compensation due to the formation of [Mn″-W<sup>•</sup>]<sup>0</sup> also might occur. The high stability of these charge complexes is confirmed by the negligible  $T_{\rm C}$  change of Mn/W-doped material with x = 0.0075 on annealing in oxidizing atmospheres at various temperatures  $T_{\rm a}$ , see Fig. 3.

## 3.2. Dielectric Constants

The dielectric constant  $\varepsilon$  and losses tan  $\delta$  of BCTZ ceramics are not only determined by the composition but also to a considerable extent by the processing conditions. In particular the atmosphere during sintering and annealing has a large influence on  $\varepsilon$  and tan  $\delta$ , as can be seen in Figs. 6(a)– 6(d) for ceramics of composition (Ba<sub>0.96</sub>Ca<sub>0.04</sub>)-[Ti<sub>0.815</sub>Zr<sub>0.18</sub>A<sub>0.005</sub>]<sub>0.995</sub>O<sub>3- $\delta$ </sub>, A = Cr, Mn, Mg, Sc. Any change of the oxygen vacancy and acceptor concentration strongly affects the dielectric properties which are closely correlated with the mobility of ferroelectric domain walls [12].

During re-oxidation of BCTZ doped with valence instable acceptors part of the oxygen vacancies are eliminated. Figure 6(a) and 6(b) shows the temperature dependence of  $\varepsilon$  and tan  $\delta$  of BCTZ containing 1 mol% Mn or Cr respectively. The solid lines represent the results obtained after sintering in reducive atmosphere and the dashed and dotted lines those obtained after annealing in oxidizing atmosphere  $(T_a = 1000^\circ \text{C}, t_a = 2 \text{ h})$  with  $\text{Po}_2 = 5 \text{ Pa}$  and  $Po_2 = 2 \cdot 10^4$  Pa (air), respectively. After re-oxidation  $\varepsilon_{\rm max}$  is considerably increased and the dielectric losses are strongly decreased in the paraelectric region. Figure 6(e) clearly shows a slight frequency dependence of  $\varepsilon$  but a large change of tan  $\delta$  for the oxidized Sc-doped material. Similar effects are usually observed in lead containing complex perovskite ceramics which are known as relaxor type dielectrics.

The change of  $\varepsilon$  and tan  $\delta$  from the reduced to the oxidized state turned out to be a reversible process. This behavior can be attributed to the disappearence of parts of  $V_0^{\star\star}$  at re-oxidation. Acceptor-oxygen vacancy complexes reduce the mobility of ferro-electric domains, thus diminishing  $\varepsilon$  and tan  $\delta$  in the ferroelectric region (pinning effect) [12]. At the high

temperature end in the paraelectric region large numbers of  $V_0^{\bullet\bullet}$  usually give rise to increased ionic conductivity [11,13,14].

However, the change of  $\varepsilon$  and  $\tan \delta$  after reoxidation also occurs for pure BCTZ, Ca-free compositions and BCTZ doped with valency stable acceptors as demonstrated in Figs. 6(c) and 6(d) for Sc- and Mg-doped BCTZ. The strong increase of  $\varepsilon$ coincides with the decrease of  $T_{\rm C}$  and appears in the range  $500^{\circ}C \le T_a \le 700^{\circ}C$ . The insulation resistance also exhibits a strong increase in this range of annealing temperatures and thus strongly suggests an elimination of oxygen vacancies. However, measurable weight changes due to elimination of  $V_0^{\bullet \bullet}$  could not be detected with TGA. The change of  $\varepsilon$ and  $\tan \delta$  might be attributed on the one hand to elimination and reorientation [17] of oxygen vacancies, but on the other hand the experimental data suggest the presence of inhomogeneous distributions of oxygen vacancies and complexes, leading to changes of the domain configuration and thus the elastic energy as indicated by different domain pattern before and after annealing, see Fig. 7.

The increase of  $\varepsilon$  and tan  $\delta$  of Sc- and Mg-doped material is more pronounced than for Mn-doped BCTZ which is due to the much smaller grain size of the Mn-doped ceramics, as outlined in the next section. For the Cr- and Mn-doped ceramics a decrease of  $\varepsilon$  is observed for lower Po<sub>2</sub> (5 Pa) and  $T_a \ge 800^{\circ}$ C which corresponds to the elimination of oxygen vacancies due to the oxidation of Cr or Mn at higher  $T_a$ . The disappearence of these effects for high Po<sub>2</sub> (in air) suggests the almost complete elimination of oxygen vacancy gradients probably present in the grains at low Po<sub>2</sub>.

3.2.1. Grain Size Effects. The change of the dielectric constant depends on the grain size affecting the stress distribution and the domain configuration in the grain. The grain size dependence of  $\varepsilon$  has been primarily investigated in ferroelectric materials far below the Curie point [14,15]. Very few has been reported in contrast about the grain size dependence of the dielectric maximum. The variation of  $\varepsilon_{\text{max}}$  with the average grain size,  $d_{\text{av}}$ , is shown in Fig. 8 for BCTZ of composition (Ba<sub>0.96</sub>Ca<sub>0.04</sub>)[Ti<sub>0.82-x</sub>Zr<sub>0.18</sub>A<sub>x</sub>]<sub>0.995</sub>O<sub>3-\delta</sub>. The data represent materials, containing various acceptor dopes in concentrations  $x \le 0.01$ . Their average grain sizes are compiled in Table 1. The circles



*Fig.* 6(a-d). Temperature dependence of the dielectric constant  $\varepsilon$  and the dielectric losses tan  $\delta$  for acceptor doped materials of composition  $(Ba_{0.96}Ca_{0.04})[Ti_{0.82-x}Zr_{0.18}A_x]_{0.995}O_{3-\delta}$  in the reduced  $(Po_2 = 2 \cdot 10^{-5} Pa$ : solid line) and oxidized  $(Po_2 = 5 Pa$ : dashed line,  $Po_2 = 2 \cdot 10^4 Pa$ : dotted line) state for (a) A = Mn, (b) A = Cr,(c) A = Sc, (d) A = Mg.



Fig. 6(e). Frequency dependence of the dielectric constant  $\varepsilon$  for oxidized (Po<sub>2</sub> = 5 Pa) acceptor doped material in the range 1 kHz to 00 kHz: A = Sc.



Fig. 7. Domain pattern for an acceptor doped BCTZ ceramic with a Curie point of  $\simeq 80^{\circ}$ C in the (a) reduced and (b) oxidized state.

represent Mn containing materials. The almost linear relationship between  $\varepsilon_{\text{max}}$  and  $d_{\text{av}}$  is most probably restricted to materials, showing a narrow distribution of grain sizes. The difference of  $\varepsilon_{\text{max}}$  between reduced and oxidized materials (Po<sub>2</sub> = 5 Pa, 2 h, 1000°C) increases with increasing  $d_{\text{av}}$  and vanishes almost at grain sizes below 2  $\mu$ m. This effect might be attributed to stress induced changes of the domain configuration. These effects are also observed for donor acceptor co-doped compositions [18].

The relatively small scatter of the data suggests that the effect of oxidizing treatment seems to be almost independent of the kind and concentration of particular acceptors. This observation is confirmed by the fact that in BCTZ the maximum dielectric constant is almost the same ( $\varepsilon_{\rm max} \simeq 16,000$ ) at a grain size of  $d_{\rm av} = 4 \ \mu {\rm m}$  for all acceptors studied, see Fig. 9. The data were calculated using the grain size dependence presented in Fig. 8.



*Fig.* 8. Maximum dielectric constant as a function of the average grain size for materials of composition  $(Ba_{0.96}Ca_{0.04})[Ti_{0.82-x}Zr_{0.18}A_x]_{0.995}O_{3-\delta}$  in the reduced (solid symbols) and the oxidized (open symbols) state for various acceptors and  $x \le 0.01$ . Circular symbols: Mn, square symbols: other acceptors.

## 4. Conclusion

The investigation of acceptor doped BCTZ has shown that the dielectric and ferroelectric properties are influenced in particular by the oxygen vacancies formed for compensation of the acceptors. Charged oxygen vacancies,  $V_0^{\bullet\bullet}$ , in combination with acceptors obviously give rise to local deformation of the perovskite unit cells, leading to local reduction of



*Fig. 9.* Concentration dependence of the maximum dielectric constant for materials of composition  $(Ba_{0.96}Ca_{0.04})$ - $[Ti_{0.82-x}Zr_{0.18}A_x]_{0.995}O_{3-\delta}$  doped with various acceptors A and  $x \le 0.01$ . The measured data were corrected to an average grain size of 4  $\mu$ m, using the grain size dependence, presented in Fig. 8.

the tetragonal c/a ratio and decrease of the Curie point  $T_{\rm C}$ . The oxygen vacancies can be eliminated by combining acceptors with donors. Elimination of  $V_0^{\star\star}$  by donor-acceptor complexes thus stabilizes the ferroelectric properties, leading to an increase of  $T_{\rm C}$ . The maximum dielectric constant reveals an almost linear increase with the average grain size. The slope depends on the Zr content and annealing conditions.

Reoxidation of acceptor doped BCTZ which has been fired before in reductive atmosphere leads to an increase of permittivity in the ferroelectric region. The effect is also pronounced in materials containing valence stable acceptors, e.g., Mg, Sc. Although in this case the number of oxygen vacancies cannot be decreased by annealing in oxidizing atmosphere, the dielectric effects suggest a reduction of the number of  $V_0^{\bullet\bullet}$ . Possibly, part of the  $V_0^{\bullet\bullet}$ -acceptor complexes pinning the ferroelectric domain walls are reorientated during re-oxidation so that the elastic energy is changed and the mobility of the domains increases.

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#### References

- 1. J.C. Jaffe, in *Piezoelectric Ceramics*, edited by J.P. Roberts and P. Popper (Academic Press, London 1971), p.237.
- T. Takada, E. Chang, and D.M. Smyth, Adv. of Ceram., 19, 147 (1987).
- 3. U. Robels and G. Arlt, J. Appl. Phys., 73, 3454 (1993).
- 4. H.-J. Hagemann and D. Hennings, J. Am. Ceram. Soc., 64, 590 (1981).
- N.H. Chan, R.K. Sharma, and D.M. Smyth, J. Am. Ceram. Soc., 65, 167 (1982).
- 6. H.-J. Hagemann, J. Phys. C, Solid State, 11, 3333 (1978).
- 7. H.-J. Hagemann and H. Ihrig, Phys. Rev. B, 20, 3871 (1979).
- D. Hennings and H. Schreinemacher, J. Europ. Ceram. Soc., 15, 795 (1995).
- 9. R.C. Kell and N.J. Helicar, Acustica, 6, 235 (1956).
- D. Hennings, A. Schnell, and G. Simon, J. Am. Ceram. Soc., 65, 539 (1982).
- D. Hennings and B. Schreinemacher, J. Europ. Ceram. Soc., 14, 463 (1994).
- 12. H. Dederichs and G. Arlt, Ferroelectrics, 68, 281 (1986).
- R. Waser, T. Baiatu, and K.H. Härdtl, J. Am. Ceram. Soc., 73, 1645 (1990).

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- R. Waser, T. Baiatu, and K.H. Härdtl, J. Am. Ceram. Soc., 73, 1663 (1990).
- 15. G. Arlt, D.F.K. Hennings, and G. de With, J. Appl. Phys., 58, 1619 (1985).
- 16. G. Arlt, Ferroelectrics, 104, 217 (1990).
- 17. U. Robels and G. Arlt, J. Appl. Phys., 73, 3454 (1993).
- P. Hansen, D. Hennings, and H. Schreinemacher, J. Am. Ceram. Soc., 81, 1369 (1998).